

Spin effects in the interaction of antiprotons with the deuteron at low and intermediate energies

Yu.N. Uzikov

**Laboratory of Nuclear Problems, Joint Institute for Nuclear
Research, Dubna, Russia**

J. Haidenbauer

**Institute for Advanced Simulation, Forschungszentrum
Juelich, Germany**

MOTIVATION

* Preparation of an intense beam of polarized antiprotons (PAX)

V. Barone et al., arXiv:hep-ex/0505054.

pbar-1H scattering in rings /F. Rathmann et al. PRL 94 (2005) 014801/

p-1H scattering at 23 MeV FILTEX, F.Rathmann et al. PRL 71 (1993) 1379
49 MeV, COSY, W. Augustinyak et al. PLB 718 (2012) 64

Other possibilities - polarized deuterium targets :

Yu.N. Uzikov and J. Haidenbauer, PRC 79 (2009) 024617;

PRC 87 (2013) 054003; PRC 88 (2013) 0227001

S. Salnikov NPA 874 (2012) 98

- or polarized ^3He target

Yu.N. U., J. Haidenbauer, B. Prmantaeyva , PRC (2011)

* Spin dependence of the elementary anti- p N amplitudes

$\bar{N} N$ interaction model developed by the Juelich Group

T. Hippchen, J. Haidenbauer, K. Holinde, V. Mull,
PRC 44, 1323 (1991)

Models A and D (differ by treatment of annihilation)

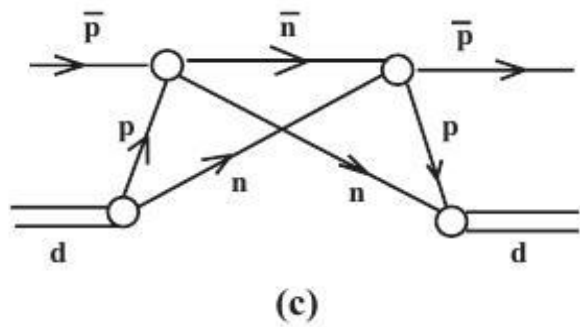
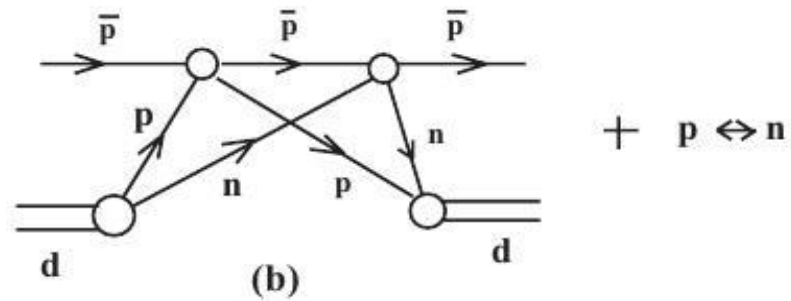
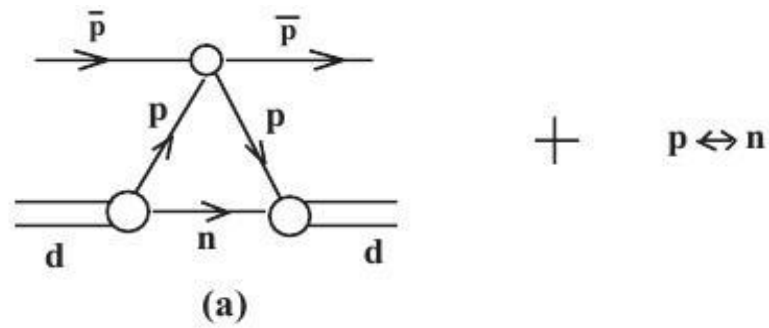
Recent partial-wave analysis of antip p - data

D. Zhou, R.G.E. Timmermans, Phys. Rev. C 86, 044003 (2012)

* \bar{p} -d dynamics -> Glauber model

pd elastic scattering: M.N. Platonova, V.I. Kukulín,
PRC 81,014004 (2010)

-> updated for \bar{p} d scattering : Yu.N. U.
J.Haidenbauer, PRC 87, 054003 (2013)



$$\hat{M}(\mathbf{q}, \mathbf{s})$$

$$= \exp\left(\frac{1}{2}i\mathbf{q} \cdot \mathbf{s}\right) M_{\bar{p}p}(\mathbf{q}) + \exp\left(-\frac{1}{2}i\mathbf{q} \cdot \mathbf{s}\right) M_{\bar{p}n}(\mathbf{q})$$

$$+ \frac{i}{2\pi^{3/2}} \int \exp(i\mathbf{q}' \cdot \mathbf{s}) [M_{\bar{p}p}(\mathbf{q}_1) M_{\bar{p}n}(\mathbf{q}_2)$$

$$+ M_{\bar{p}n}(\mathbf{q}_1) M_{\bar{p}p}(\mathbf{q}_2) - M_{\bar{p}p \rightarrow \bar{n}n}(\mathbf{q}_1) M_{\bar{n}n \rightarrow \bar{p}p}(\mathbf{q}_2)] d^2\mathbf{q}'.$$

$$M_{\bar{p}N} = A_N + (C_N \boldsymbol{\sigma}_1 + C'_N \boldsymbol{\sigma}_2) \cdot \hat{\mathbf{n}} + B_N (\boldsymbol{\sigma}_1 \cdot \hat{\mathbf{k}}) (\boldsymbol{\sigma}_2 \cdot \hat{\mathbf{k}}) \\ + (G_N - H_N) (\boldsymbol{\sigma}_1 \cdot \hat{\mathbf{n}}) (\boldsymbol{\sigma}_2 \cdot \hat{\mathbf{n}}) + (G_N + H_N) (\boldsymbol{\sigma}_1 \cdot \hat{\mathbf{q}}) (\boldsymbol{\sigma}_2 \cdot \hat{\mathbf{q}})$$

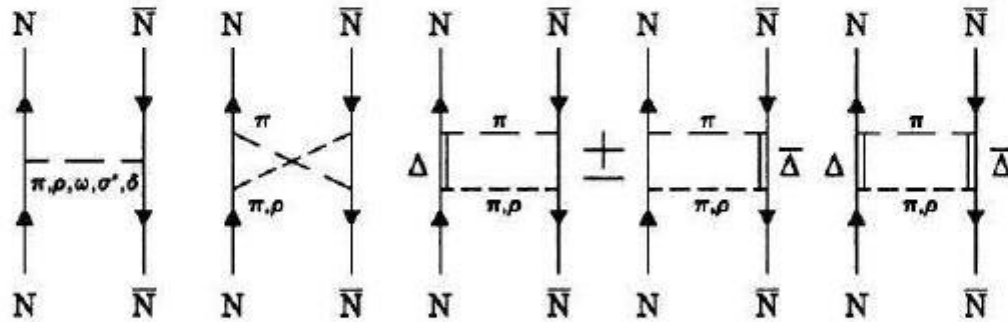
The unit vectors are defined by $\hat{\mathbf{k}} = (\mathbf{k}_i + \mathbf{k}_f)/|\mathbf{k}_i + \mathbf{k}_f|$, $\hat{\mathbf{q}} = (\mathbf{k}_i - \mathbf{k}_f)/|\mathbf{k}_i - \mathbf{k}_f|$, and $\hat{\mathbf{n}} = [\hat{\mathbf{k}} \times \hat{\mathbf{q}}]$, where \mathbf{k}_i (\mathbf{k}_f) denotes the momentum of the incident (outgoing)

The Jülich NN model

I) V_{el}

starting point: Bonn NN potential

(R. Machleidt, K. Holinde, C. Elster, Phys. Rep. 149 (1986) 1)



G-parity transform

well defined over whole range

no modification of short-range part is required

The Jülich NN model

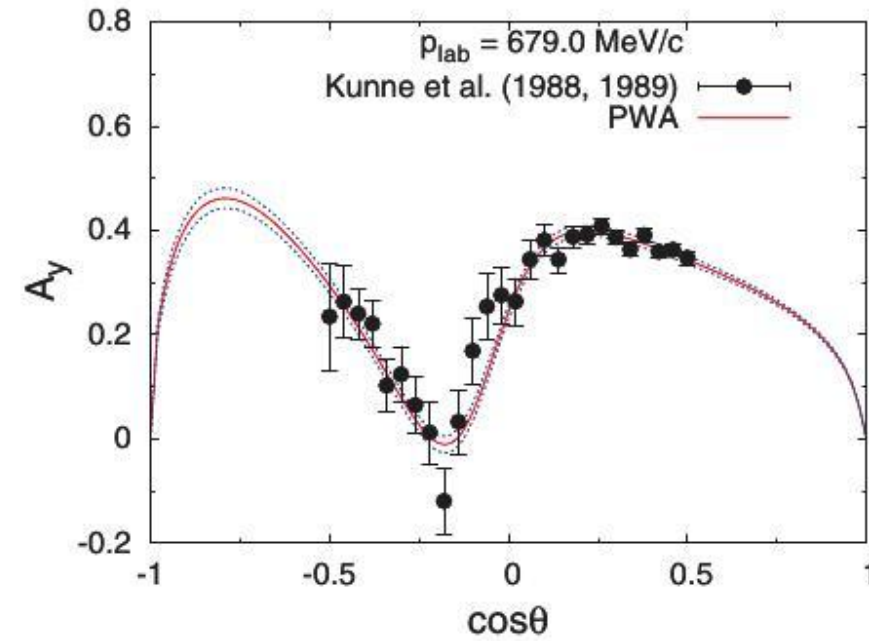
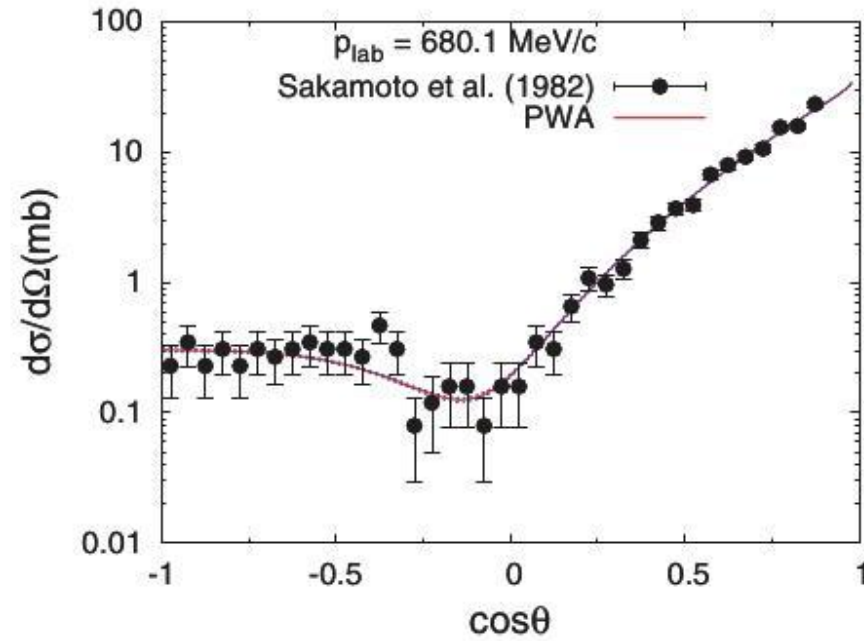
- microscopic annihilation model (for 2-meson channels) (D)



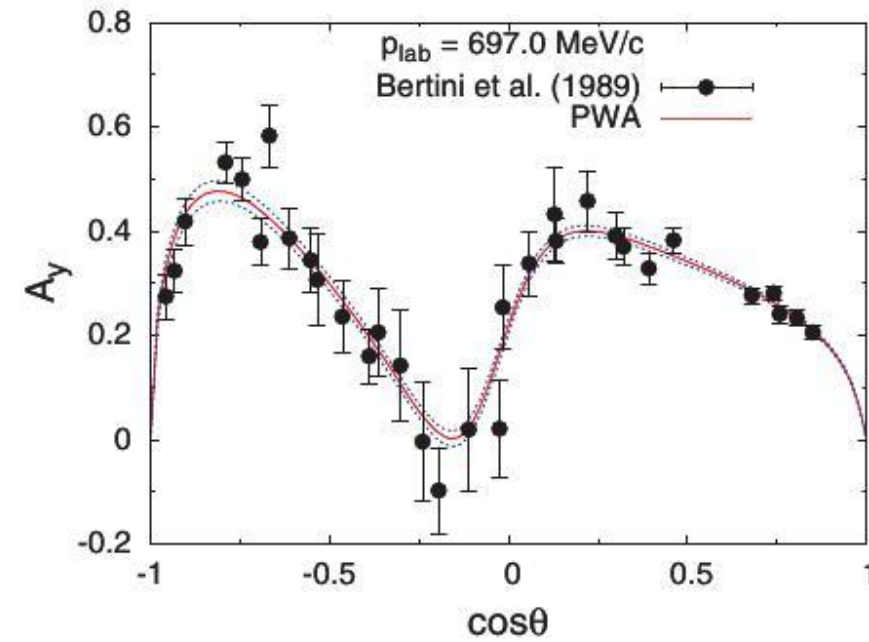
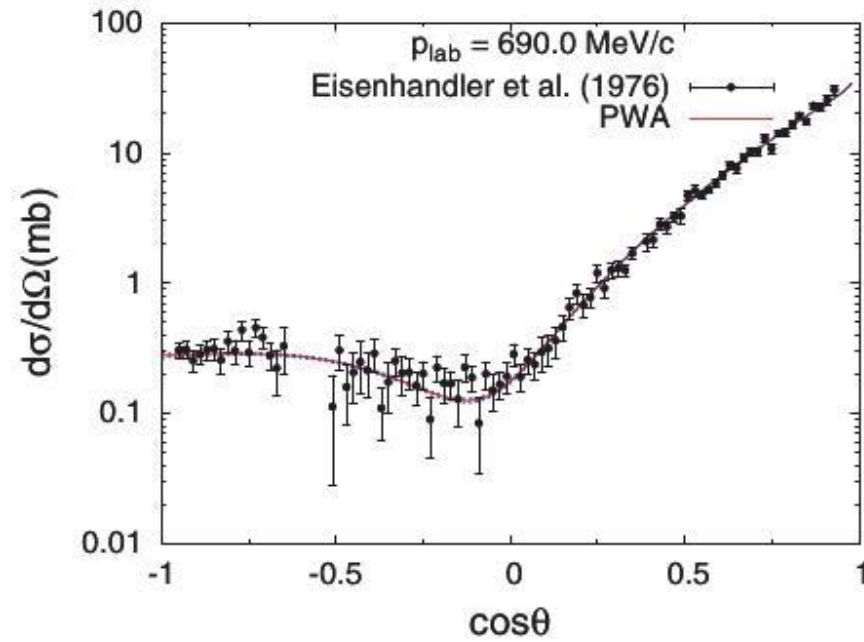
$$V^{N\bar{N} \rightarrow M_i M_j} = \begin{array}{c} M_i \quad M_j \quad \pi, \rho \quad \pi, \rho \quad K, K^* \quad K, K^* \\ \vdots \quad \vdots \quad \vdots \quad \vdots \quad \vdots \quad \vdots \\ \text{---} \quad \text{---} \quad \text{---} \quad \text{---} \quad \text{---} \quad \text{---} \\ \text{N} \quad \text{N} \quad \Delta \quad \text{N} \quad \text{N} \quad \text{N} \\ \uparrow \quad \downarrow \quad \uparrow \quad \downarrow \quad \uparrow \quad \downarrow \\ \text{N} \quad \bar{N} \quad \text{N} \quad \bar{N} \quad \text{N} \quad \bar{N} \end{array}$$

$$M_{i,j} = \pi, \eta, \rho, \omega, f_0, a_0, f_1, a_1, f_2, a_2$$

- T. Hippchen et al., PRC 44 (1991) 1323; V. Mull et al., PRC 44 (1991) 1337
- V. Mull & K. Holinde, PRC 51 (1995) 2360



ZT



Optical theorem:

$$\text{Im} \frac{\text{Tr}(\hat{\rho}_i \hat{F}(0))}{\text{Tr} \hat{\rho}_i} = \frac{k_{\bar{p}d}}{4\pi} \sigma_i.$$

M.P. Rekaló et al. Few -Bod. Syst (1998):

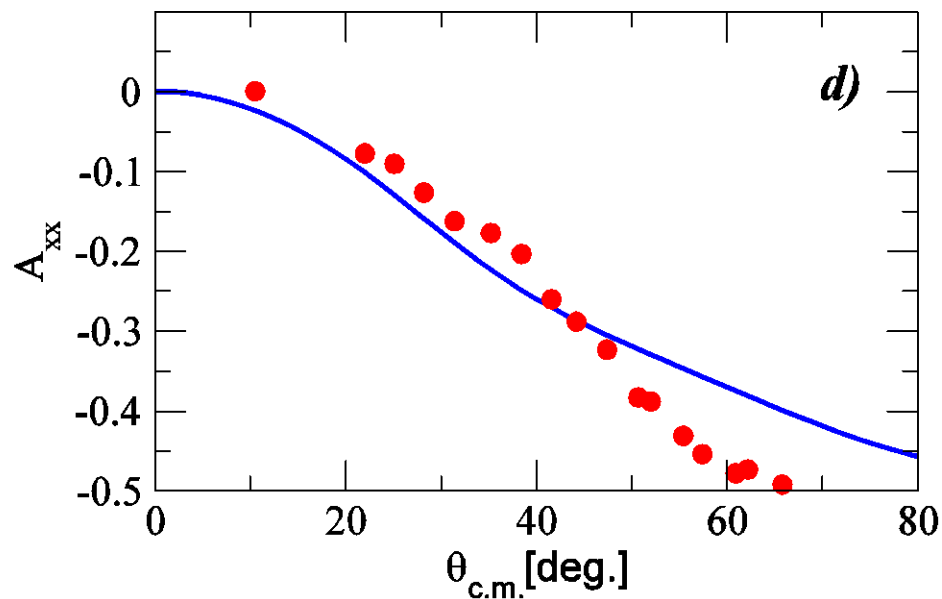
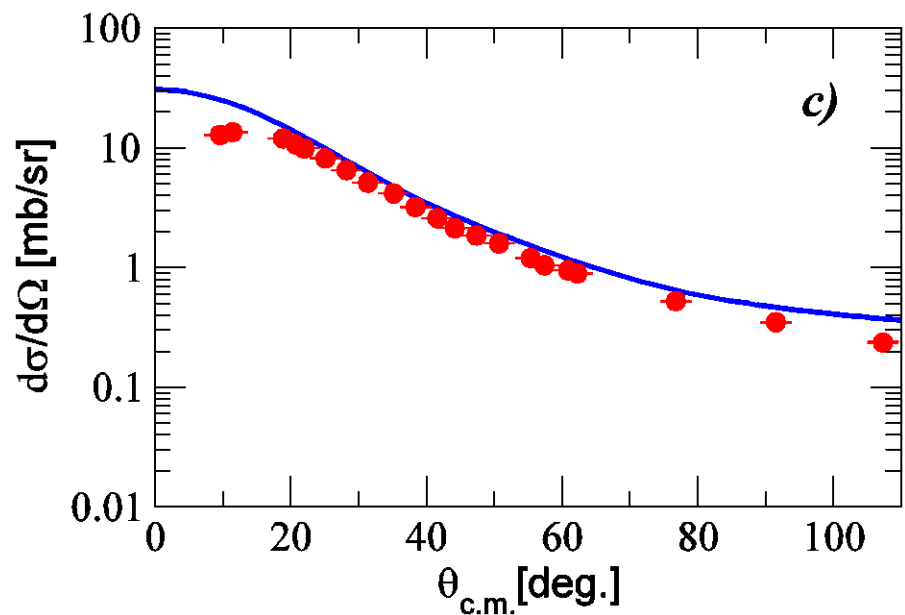
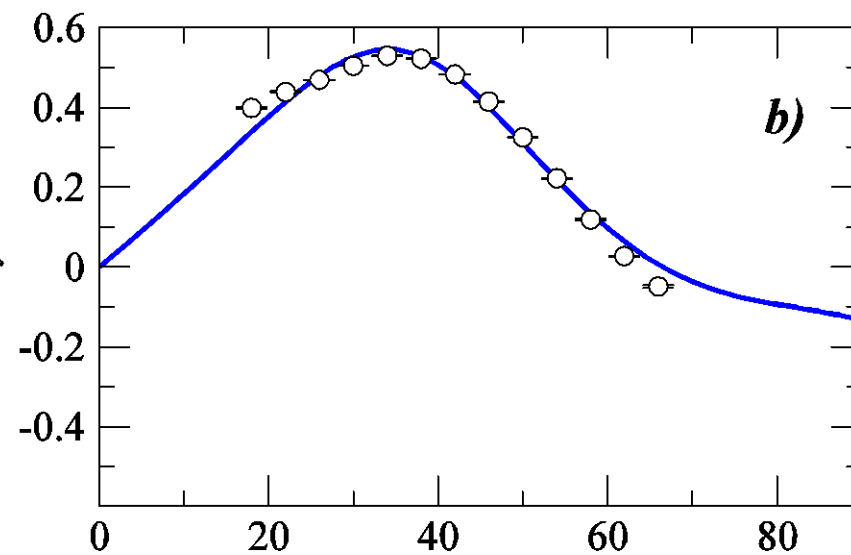
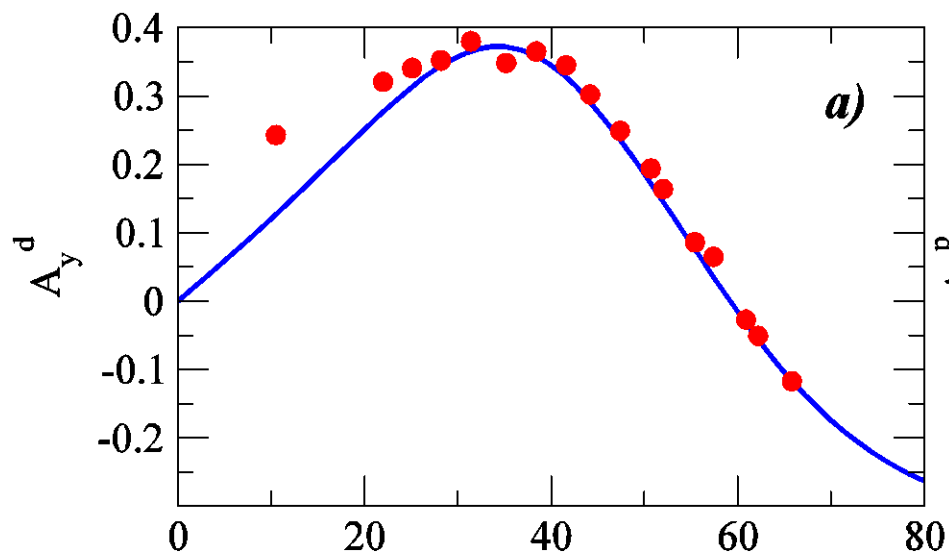
$$\begin{aligned} \hat{F}_{\alpha\beta}(0) = & g_1 \delta_{\alpha\beta} + (g_2 - g_1) m_\alpha m_\beta + i g_3 \hat{\sigma}_i \epsilon_{\alpha\beta i} \\ & + i(g_4 - g_3) \hat{\sigma}_i m_i m_j \epsilon_{\alpha\beta j}, \end{aligned}$$

The total $\bar{p}d$ cross section is defined by 11

$$\sigma = \sigma_0 + \sigma_1 \mathbf{P}^{\bar{p}} \cdot \mathbf{P}^d + \sigma_2 (\mathbf{P}^{\bar{p}} \cdot \hat{\mathbf{k}})(\mathbf{P}^d \cdot \hat{\mathbf{k}}) + \sigma_3 P_{zz}, \quad (1)$$

TEST calculation: pd elastic scattering at 135 MeV (Yu.N. U., J. Haidenbauer , arxiv 1212.2761)

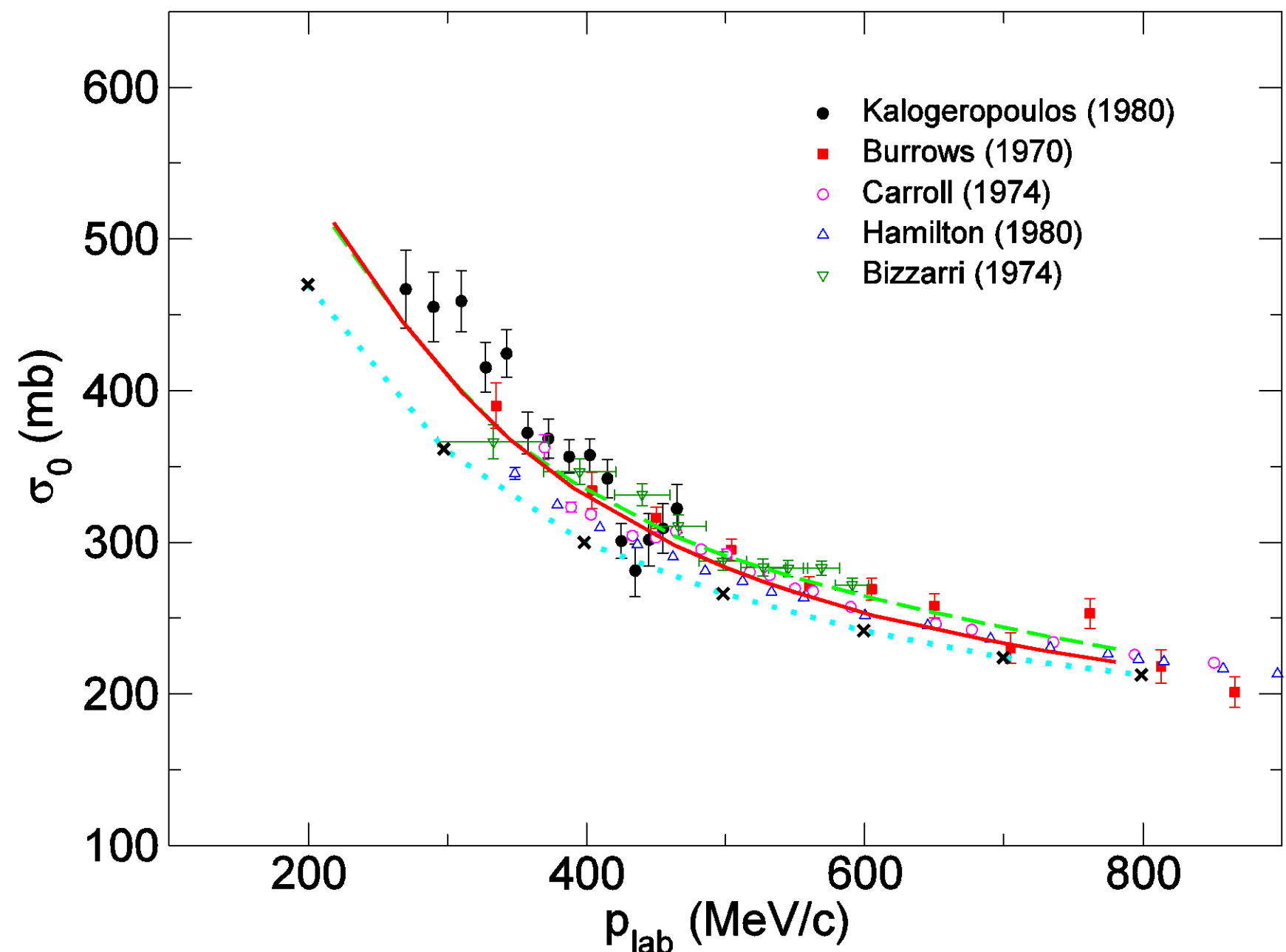
data :K. Sekiguchi et al. (2002) ; B. von Przewoski et al (2006)

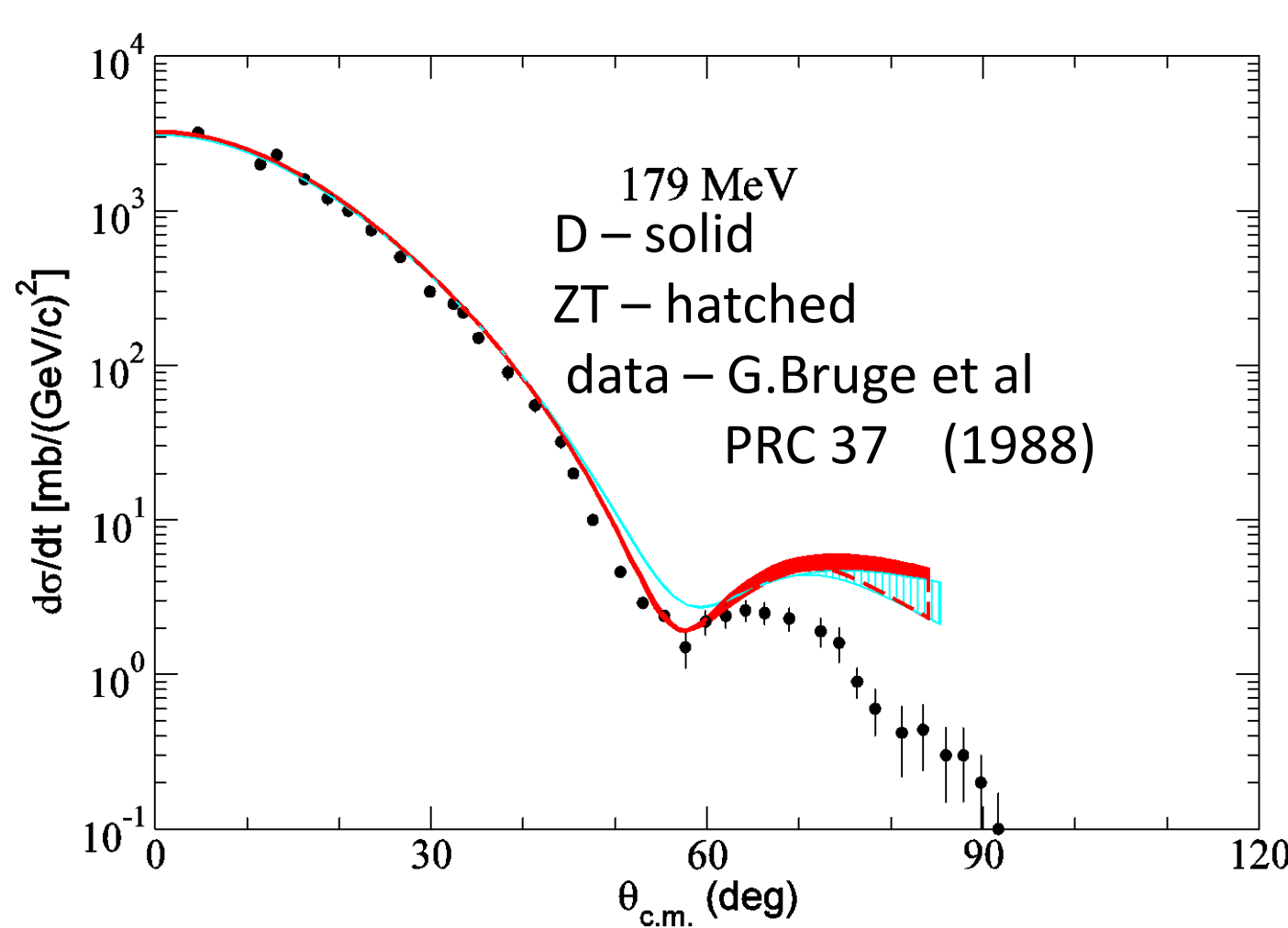
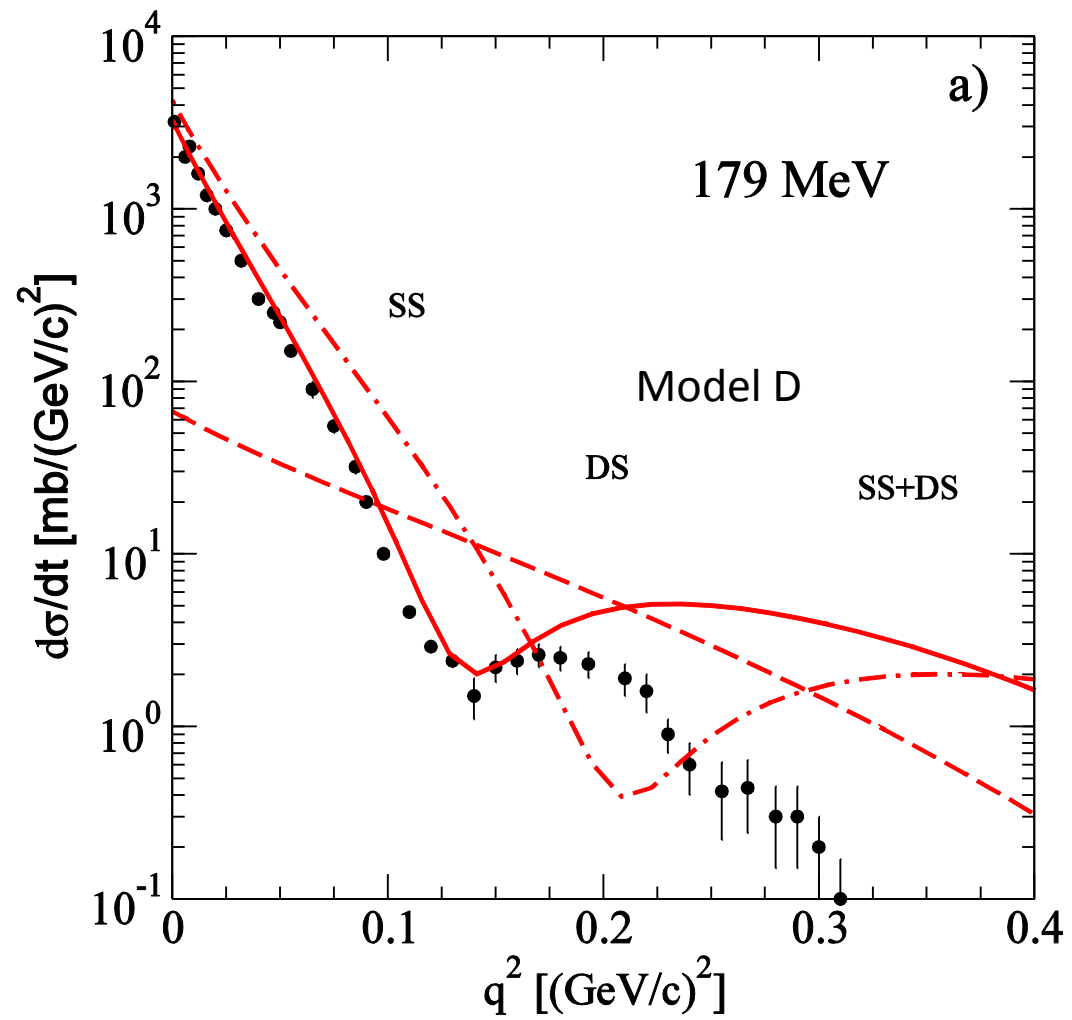


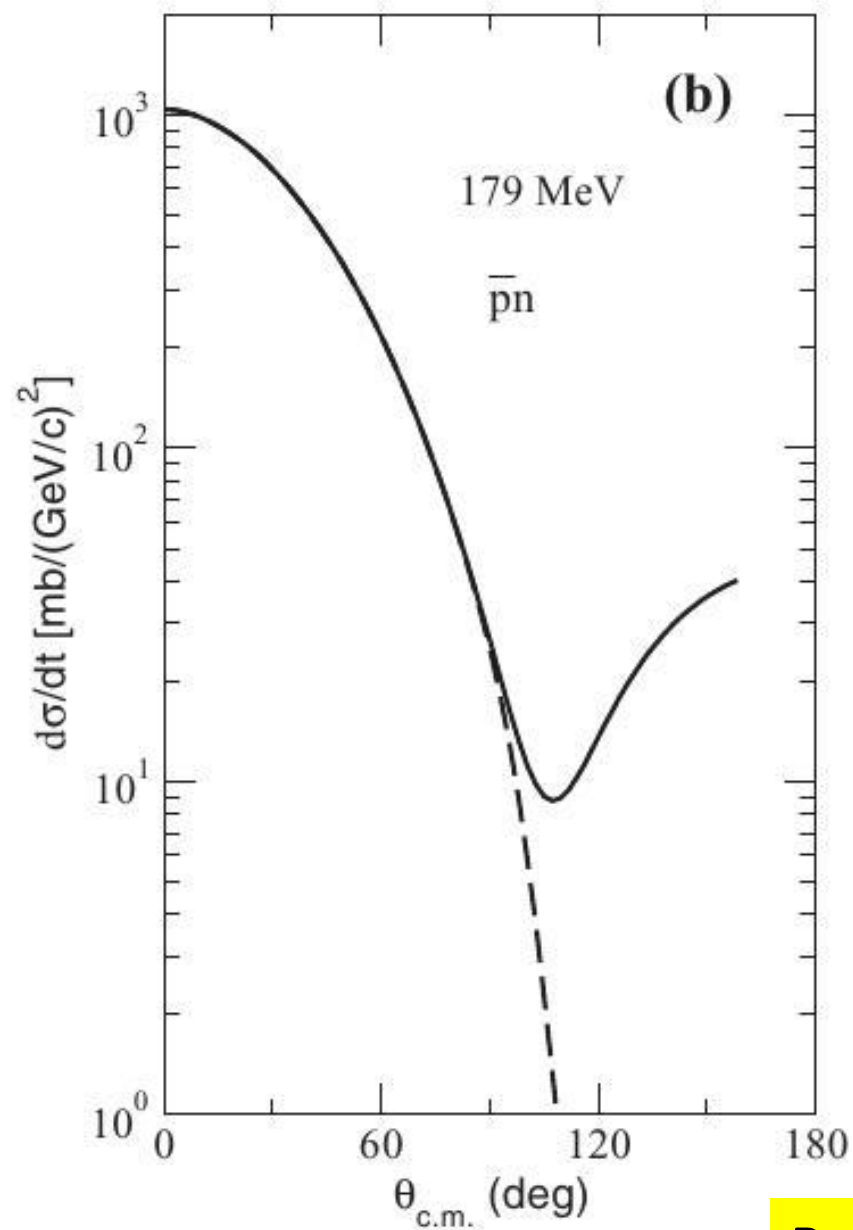
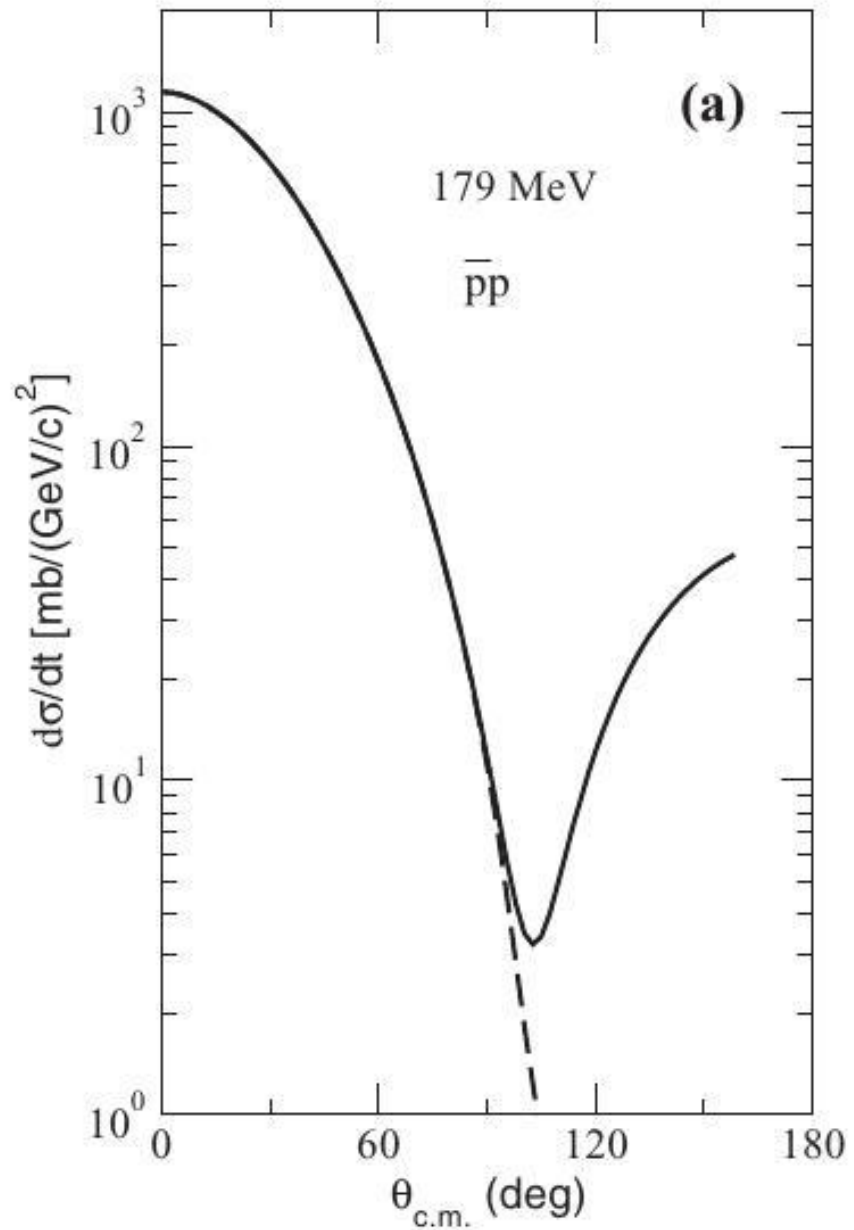
Total unpolarized $p\bar{p}$ -d cross section

Yu.N.Uzikov, J. Haidenbauer, PRC 88 (2013) 027001

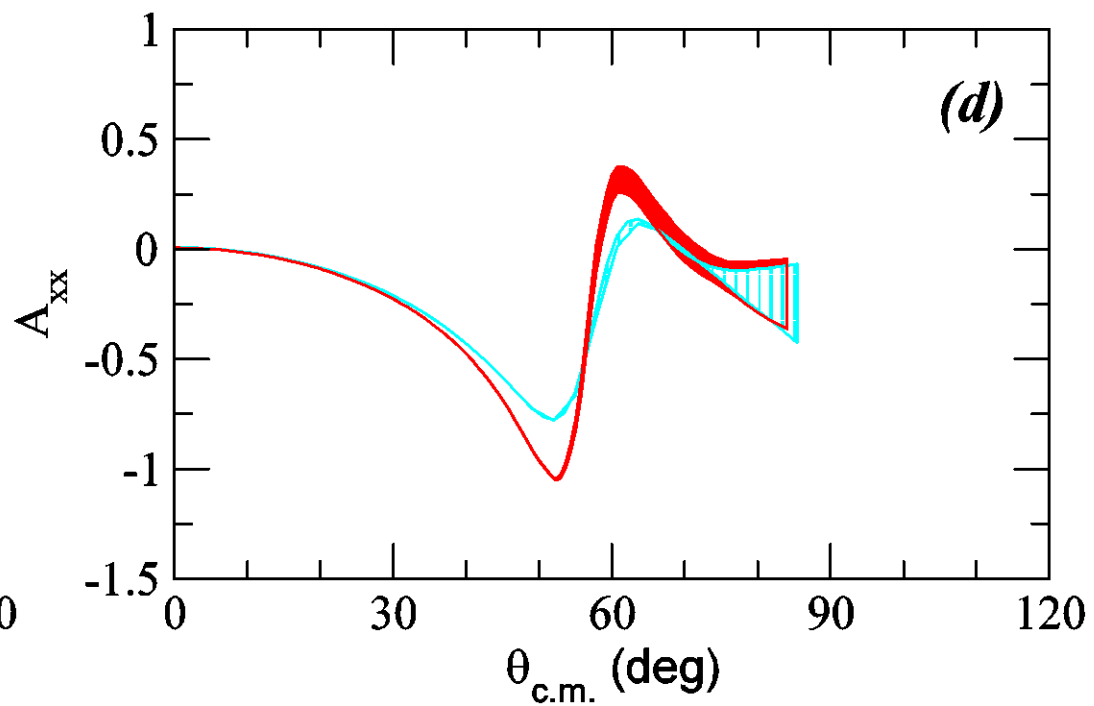
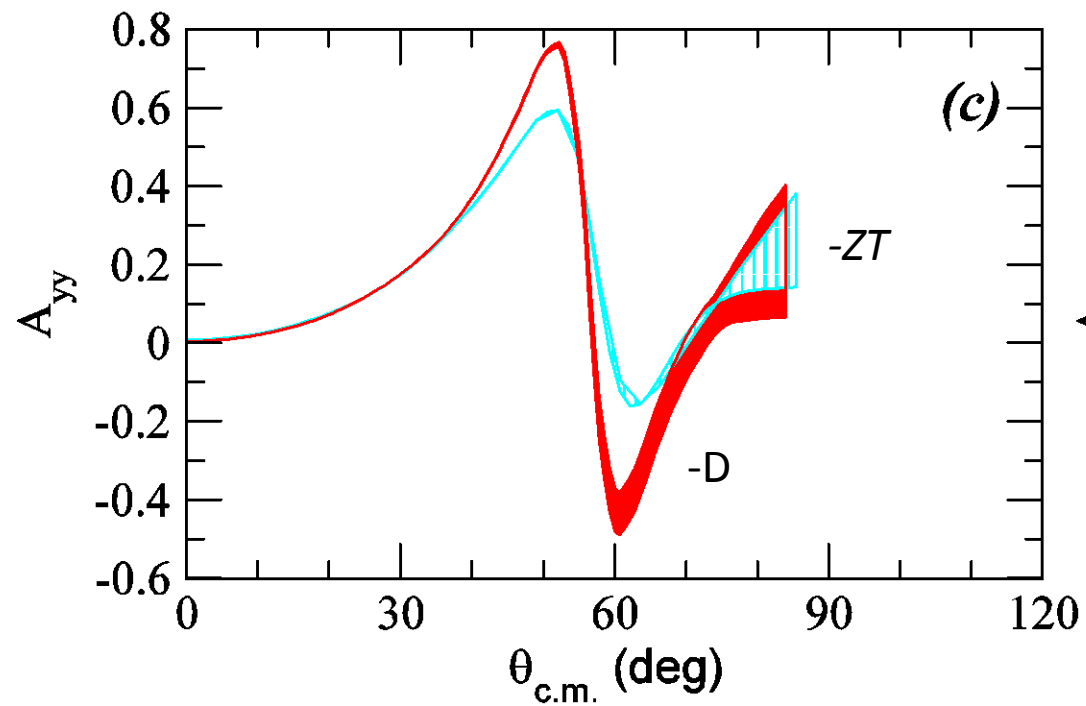
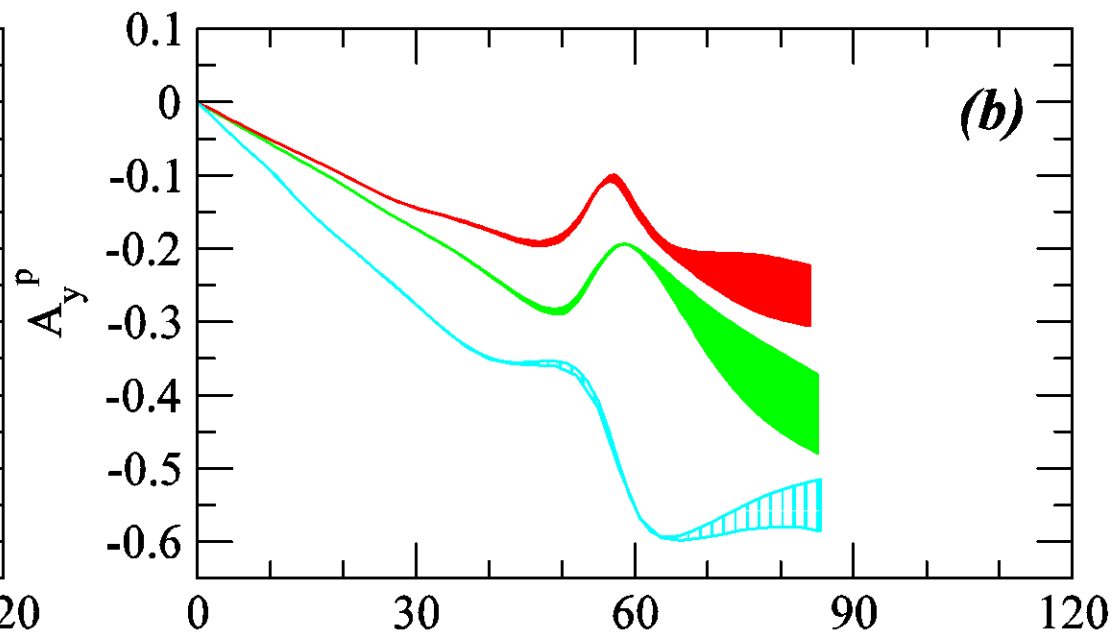
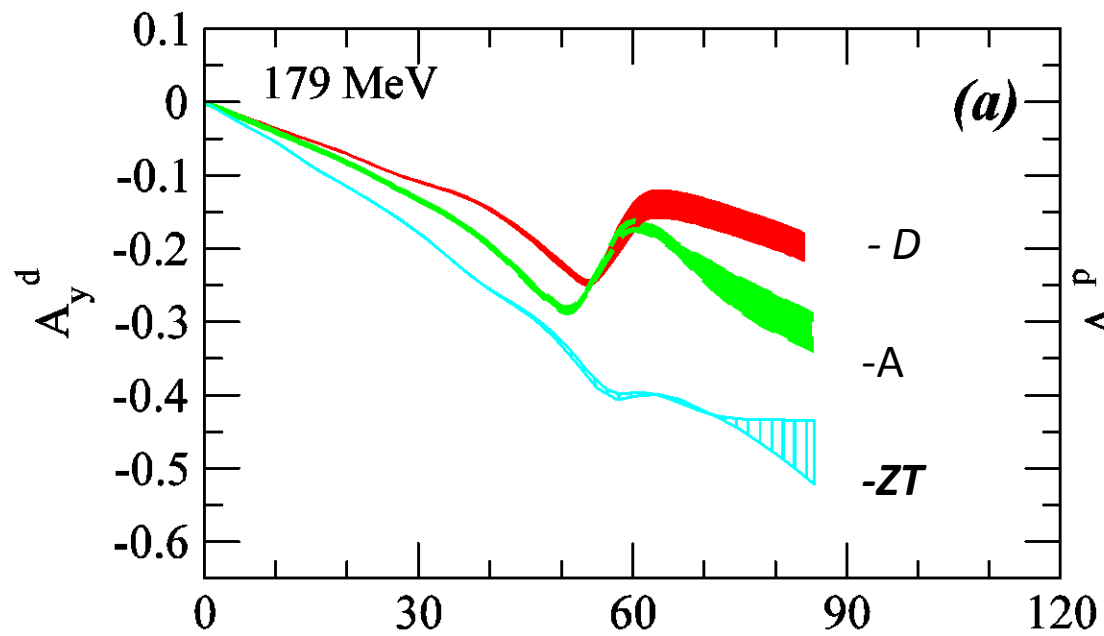
dotted – ZT
red – D
green – A

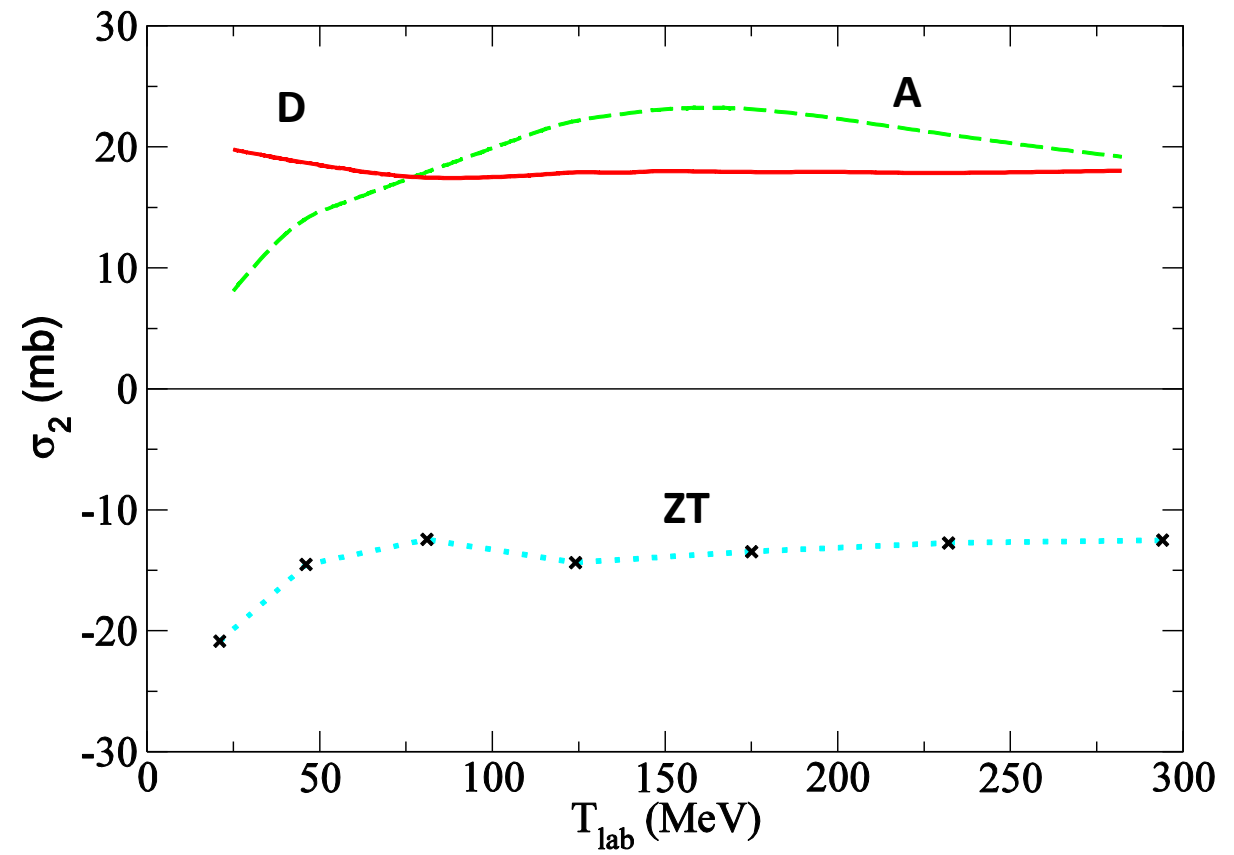
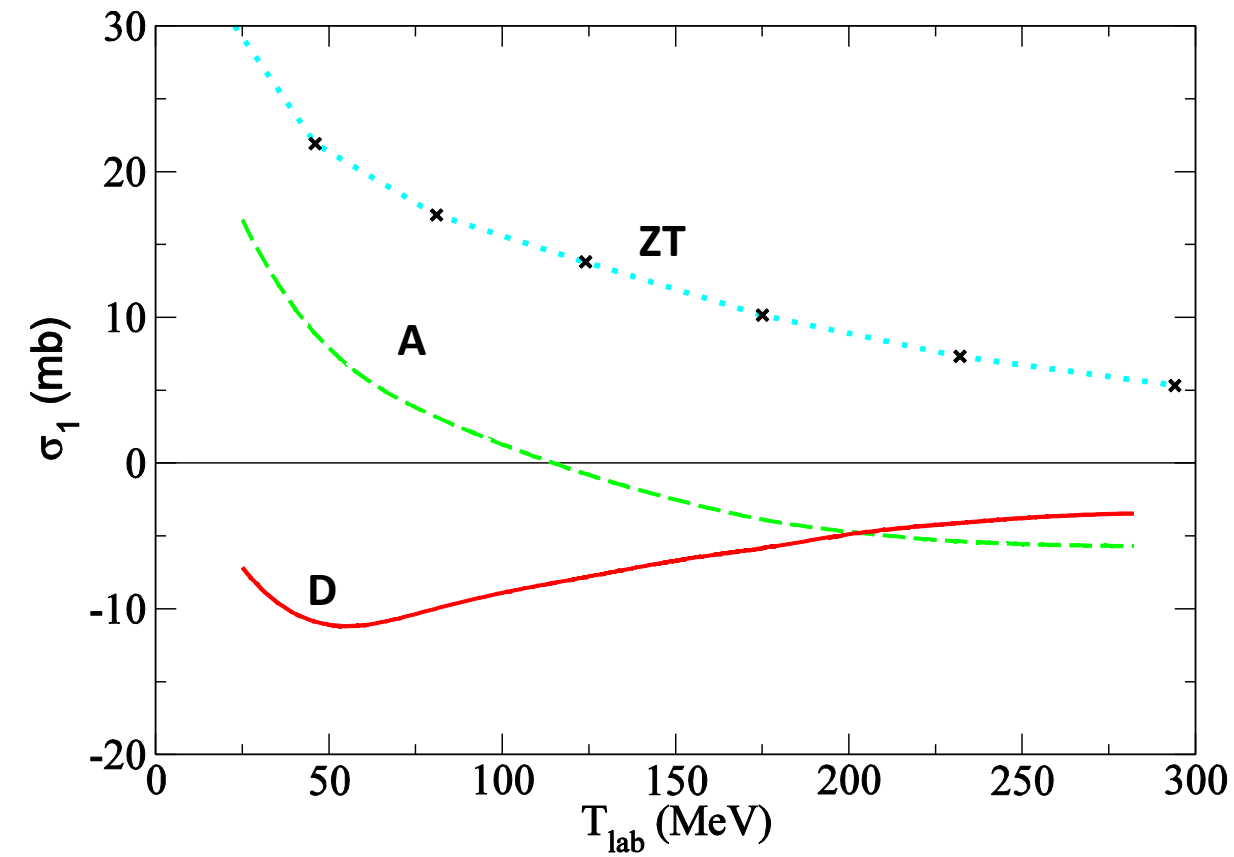






Backward tail and cut off

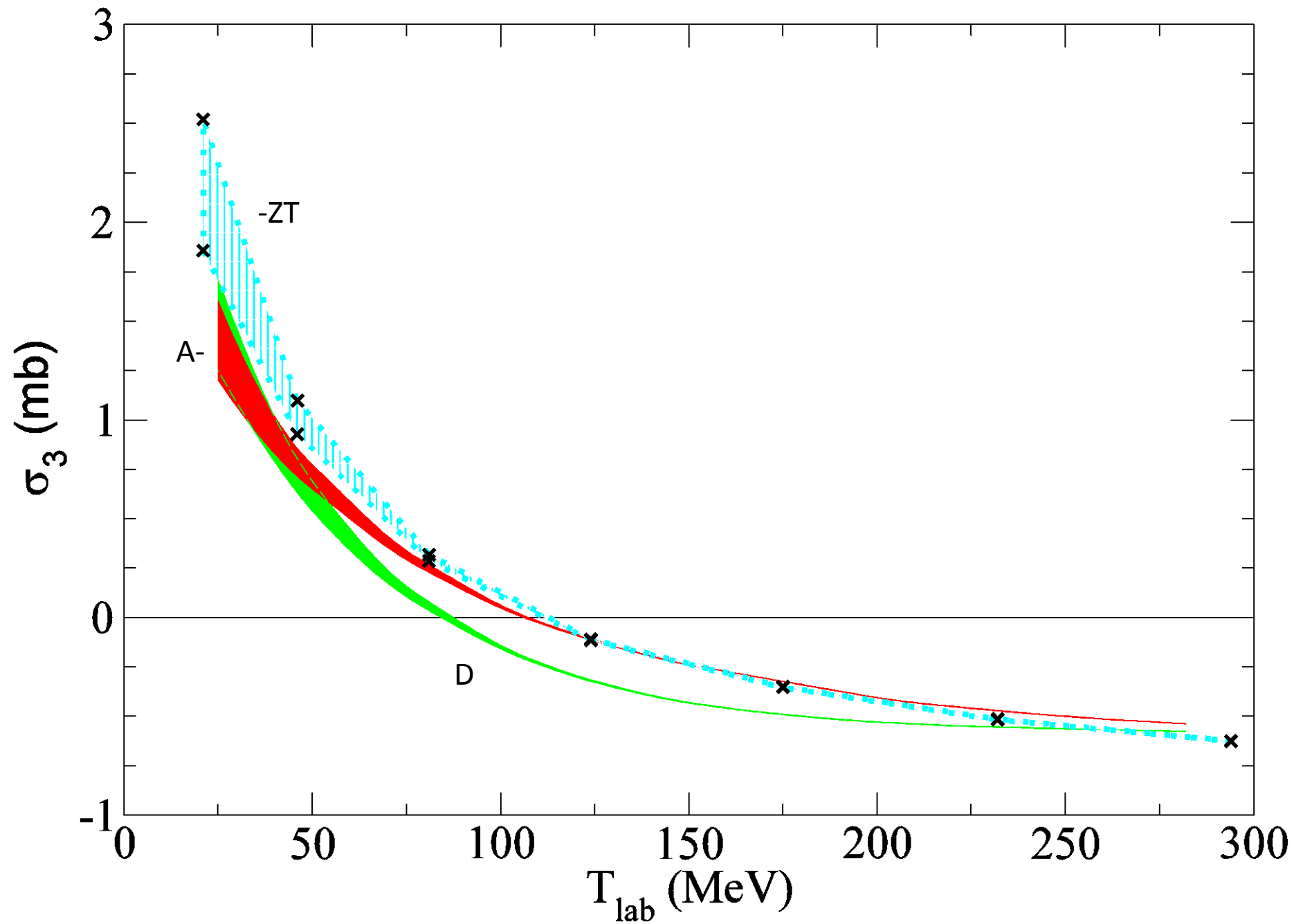


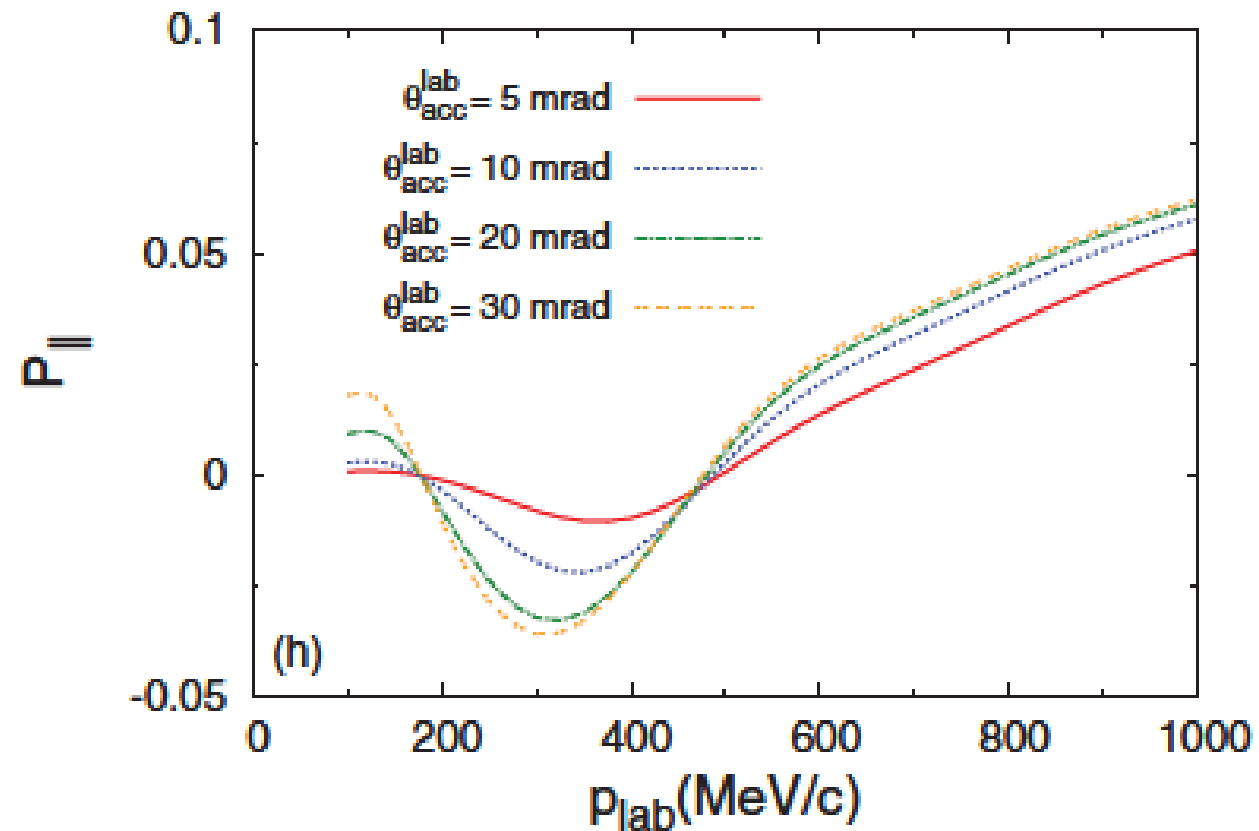
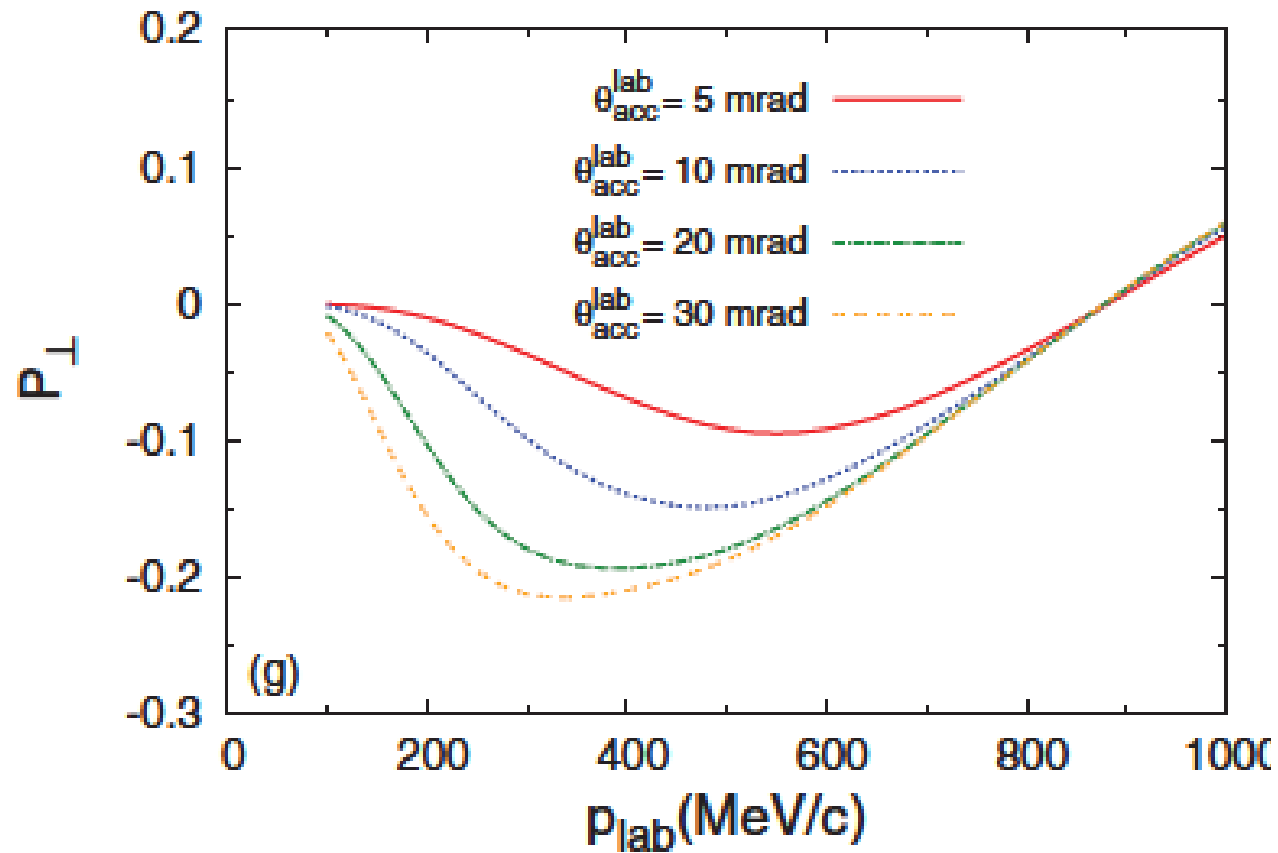
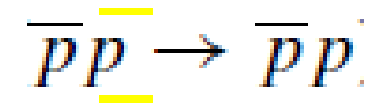


$$P_{\bar{p}}(t_0) = -2P_T \frac{\sigma_1}{\sigma_0}, \text{ if } \zeta \cdot \hat{\mathbf{k}} = 0,$$

$$P_{\bar{p}}(t_0) = -2P_T \frac{\sigma_1 + \sigma_2}{\sigma_0}, \text{ if } |\zeta \cdot \hat{\mathbf{k}}| = 1, \quad (2)$$

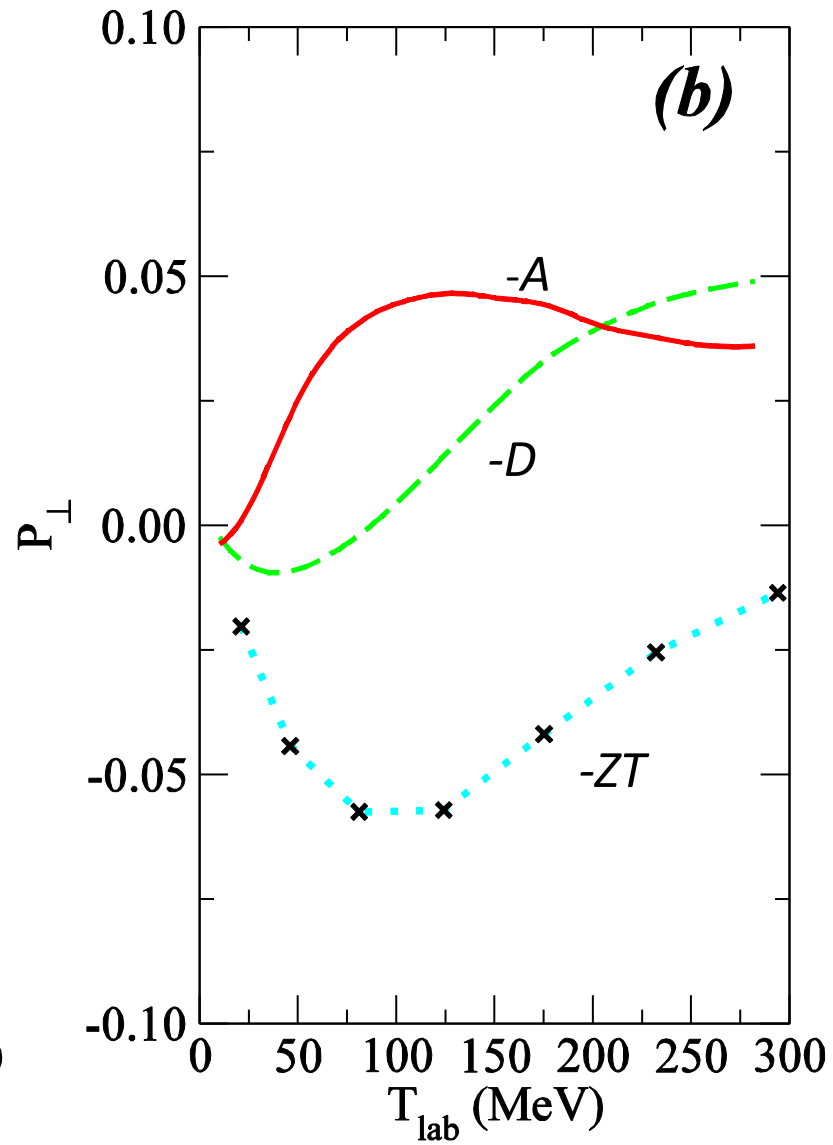
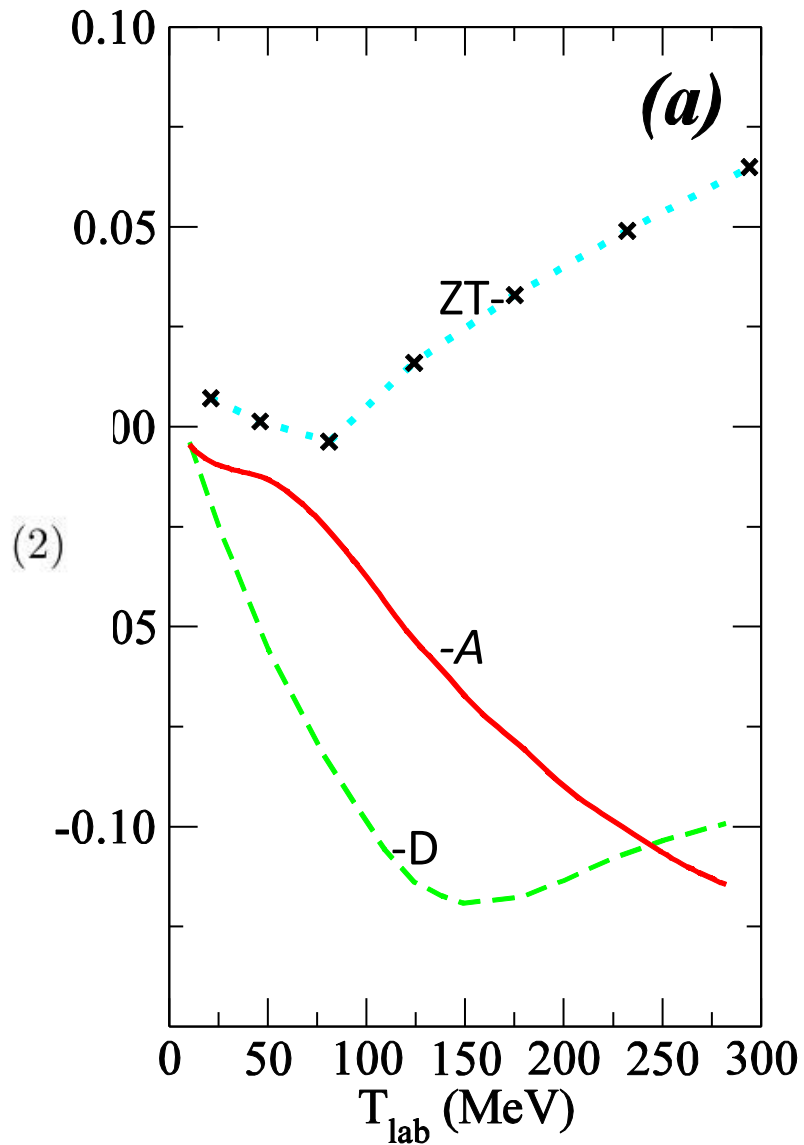
V.F. Dmitriev, A. Milstein, S. Salnikov, PLB 690 (2010) 427





$$P_{\bar{p}}(t_0) = -2P_T \frac{\sigma_1}{\sigma_0}, \text{ if } \zeta \cdot \hat{\mathbf{k}} = 0,$$

$$P_{\bar{p}}(t_0) = -2P_T \frac{\sigma_1 + \sigma_2}{\sigma_0}, \text{ if } |\zeta \cdot \hat{\mathbf{k}}| = 1, \quad (2)$$



SUMMARY

- Glauber theory is applied to \bar{p} -d scattering accounting for SS and DS mechanisms
- Full spin structure of elementary \bar{p} -N amplitudes is included
- Amplitudes from the Juelich $\bar{N}N$ model and from recent model-independent analysis by Zhou-Timmermans are employed
- Integrated cross sections and spin observables A_y , A_{xx} , A_{yy} are calculated
- Polarization efficiency for antiprotons in rings is calculated at 50-300 MeV
- Agreement with measured total cross section and with diff. cross section at 179 MeV
- Polarization efficiency for a deuterium target is comparable to 1H target
- Need more \bar{p} d data ! AD at CERN or at FAIR in Darmstadt ?